

## Real-space observation of ligand holes in SrFeO<sub>3</sub>

An anomalously high valence state of a transition-metal element sometimes shows up in oxide compounds. In such systems, holes tend to occupy mainly the ligand *p* orbitals, giving rise to interesting physical properties such as superconductivity in cuprates and rich magnetic phases in ferrates.

SrFeO<sub>3</sub> is an archetypal tetravalent ferrate compound, in which, at least formally, four electrons occupy the Fe 3*d* orbitals. It forms the perovskite-type structure with the cubic space group *Pm* $\bar{3}$ *m* (Fig. 1(a)) and exhibits metallic conductivity [1]. The local symmetry at the Fe site is *m* $\bar{3}$ *m*. Each Fe atom is surrounded by an O<sub>6</sub> regular octahedron without Jahn-Teller distortion, where the 3*d* orbitals are split into the lower-lying triplet (*t*<sub>2*g*</sub>) and the higher-lying doublet (*e*<sub>g</sub>). The high-spin state of 3*d*<sup>4</sup> corresponds to the *t*<sub>2*g*</sub><sup>3</sup>*e*<sub>g</sub><sup>1</sup> electron configuration, which causes some anisotropy in the valence electron density. However, previous X-ray photoelectron spectroscopy and X-ray absorption spectroscopy measurements suggest that the ground state consists of mixed 3*d*<sup>4</sup> and 3*d*<sup>5</sup> $\underline{L}$  ( $\underline{L}$ : ligand hole) states [2,3]. In the extreme limit of the 3*d*<sup>5</sup> $\underline{L}$  state, the electron density around the Fe site should be spherical because of the *t*<sub>2*g*</sub><sup>3</sup>*e*<sub>g</sub><sup>2</sup> electron configuration.

Although the ligand holes in the crystal have been observed by spectroscopy measurements [2,3], no one has ever seen where the ligand holes exist in real space. To observe the spatial distribution of the holes, the measurement with high-wavevector (*Q*) resolution is indispensable. In this study, we observe the valence electron density distribution of SrFeO<sub>3</sub> by electron density analysis using state-of-the-art synchrotron X-ray diffraction (XRD).

XRD experiments were performed at SPring-8 BL02B1. A He-gas-blowing device was employed to cool the crystal to 30 K. The X-ray wavelength was  $\lambda = 0.31020$  Å. A two-dimensional detector CdTe PILATUS, which had a dynamic range of  $\sim 10^7$ , was used to record the diffraction pattern. The intensities of Bragg reflections with the interplane distance  $d > 0.28$  Å were collected. A core differential Fourier synthesis (CDFS) method [4] was used to extract the valence electron density distribution around each atomic site. [Kr], [Ar], and [He] type electron configurations were regarded as core electrons for Sr, Fe, and O atoms, respectively.

Figure 1(b) shows the valence electron density distribution of SrFeO<sub>3</sub> at 30 K. No valence electron density larger than  $3e/\text{Å}^3$  is observed around the Sr site, which is consistent with the Sr<sup>2+</sup> (5*s*<sup>0</sup>) state. In contrast, valence electrons are observed around the Fe and O sites, as shown by yellow iso-density surfaces. An orange iso-density surface for higher electron distribution is observed only around the Fe site, which is clearly distinct from a sphere: there are six hollow holes toward the six ligand oxygens.

To quantify the anisotropy of the valence electron density  $\rho(\mathbf{r})$  around the Fe site, the density at a distance  $r = 0.2$  Å from the Fe nucleus is shown by a color map on a sphere (Fig. 2(a)). The maximum and minimum electron densities are present along the  $\langle 111 \rangle$  and  $\langle 100 \rangle$  axes, respectively. Figures 2(b) and 2(c) show surface color maps of  $\rho(\theta, \phi)$  for the calculated electron density considering the high-spin 3*d*<sup>4</sup> and 3*d*<sup>5</sup> states for an isolated Fe ion, respectively. In the case of 3*d*<sup>4</sup>, we assume that an

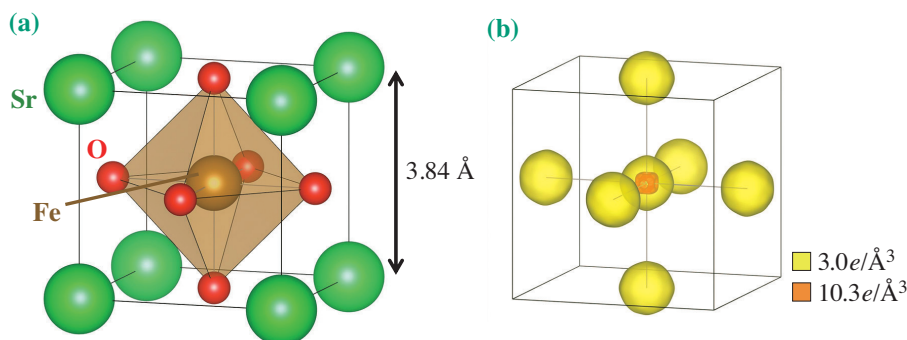


Fig. 1. (a) Crystal structure and (b) valence electron density distribution of SrFeO<sub>3</sub> at 30 K. Yellow and orange iso-density surfaces show electron-density levels of  $3.0e/\text{Å}^3$  and  $10.3e/\text{Å}^3$ , respectively.

electron occupies each  $e_g$  orbital with a probability of  $1/2$ . A clear anisotropy is observed in the  $3d^4$  state, in contrast to the completely isotropic electron density in the  $3d^5$  state. By comparing the CDFS results and simulations, the number of Fe  $3d$  electrons is estimated to be  $N_e = 4.64(8)$ , which is consistent with the previous reports of X-ray absorption spectroscopy measurement ( $N_e = 4.7$ ) [3].

Since the corresponding valence of Fe obtained by the CDFS analysis was 3.36(8), the oxygen valence is estimated to be  $-1.79(3)$ , which deviates from the ideal closed-shell value of  $-2$ . That is, the valence electron density distribution around the O site should

not be isotropic. Figure 2(d) shows a color map of the electron density  $\rho(\theta, \phi)$  at a distance  $r = 0.40 \text{ \AA}$  from the O nucleus, obtained from the CDFS analysis. The observed  $\rho(\theta, \phi)$  has some anisotropy, which differs from the isotropic behavior of an ideal  $O^{2-}$  ion (Fig. 2(e)). While the highest electron density exists toward the surrounding four Sr atoms, the lowest electron density is observed in the  $[100]$  direction toward Fe. These results suggest the existence of ligand holes accommodated in the  $O 2p_\sigma - Fe 3d$  antibonding  $\sigma^*$  orbital. The distribution of ligand holes around the O site was captured for the first time by the CDFS analysis [5].

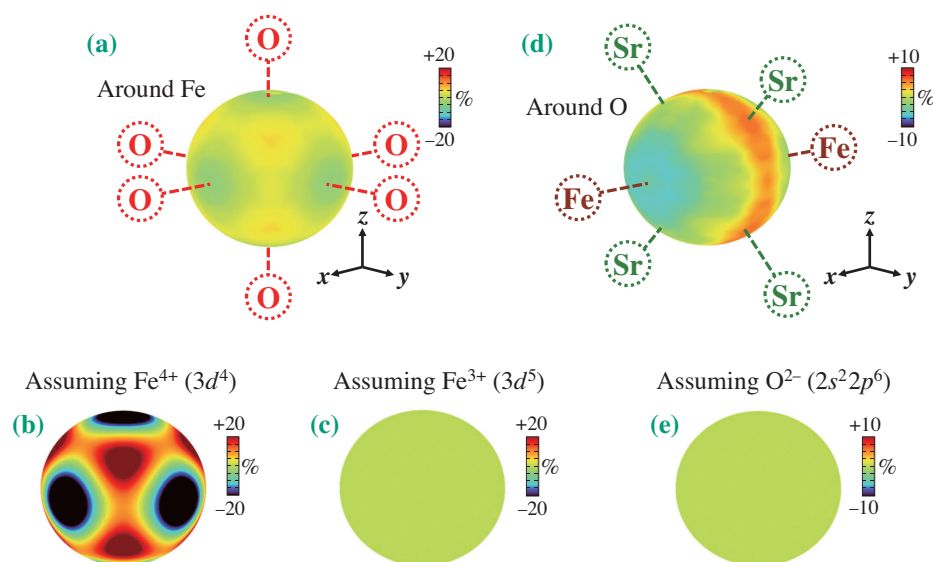


Fig. 2. Color map of the electron density (a) at a distance  $r = 0.2 \text{ \AA}$  from the Fe nucleus and (b) at a distance  $r = 0.4 \text{ \AA}$  from the O nucleus. Color maps of the calculated direction dependence of electron density for the (c)  $Fe^{4+} 3d^4$ , (d)  $Fe^{3+} 3d^5$  and (e)  $O^{2-} 2s^2 2p^6$  states assuming isolated Fe and O atoms. The color bar scale is represented by  $\frac{[\rho(\theta, \phi) - N_e/4\pi]}{N_e/4\pi} \times 100 (\%)$ .  $N_e$  is the number of valence electrons.

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## References

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